

$$(a) \quad \dot{m} c \, dT = U P \, dx (T_a - T)$$

$$\frac{dT}{T_a - T} = \frac{UP}{\dot{m}c} \, dx$$

$$\ln \frac{(T_{ai} - T_{fi})}{(T_{ai} - T_{si})} = - \frac{UP}{\dot{m}c} L$$

$$T_{fi} = T_{ai} - (T_{ai} - T_{si}) \exp \left(- \frac{UP}{\dot{m}c} L \right)$$

$$(b) \quad \epsilon = \frac{Q}{Q_{MAX}}$$

$$\begin{aligned} Q_i &= \dot{m} c (T_{fi} - T_{si}) = \dot{m} c \left[T_{ai} - T_{si} - (T_{ai} - T_{si}) \exp \left(- \frac{UP}{\dot{m}c} L \right) \right] \\ &= \dot{m} c (T_{ai} - T_{si}) \left[1 - \exp \left(- \frac{UP}{\dot{m}c} L \right) \right] \end{aligned}$$

$$Q_{MAX} = \dot{m} c (T_{ai} - T_{si})$$

$$\epsilon = \frac{Q_i}{Q_{MAX}} = \left[1 - \exp \left(- \frac{UP}{\dot{m}c} L \right) \right]$$

$$(c) \quad Q_i = \dot{m} c (T_{ai} - T_{si}) \epsilon$$

$$Q_{i+1} = \dot{m} c (T_{ai+1} - T_{si+1}) \epsilon$$

$$\frac{Q_{i+1}}{Q_i} = \frac{(T_{ai+1} - T_{si+1})}{(T_{ai} - T_{si})}$$

$$\dot{m}_2 c_2 (T_{ai} - T_{ai+1}) = \dot{m} c (T_{ai} - T_{si}) \epsilon$$

$$T_{ai+1} = T_{ai} - \frac{\dot{m} c}{\dot{m}_2 c_2} (T_{ai} - T_{si}) \epsilon$$

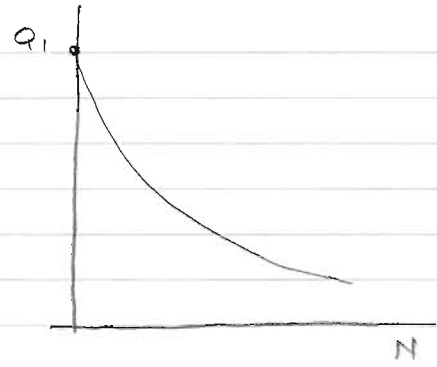
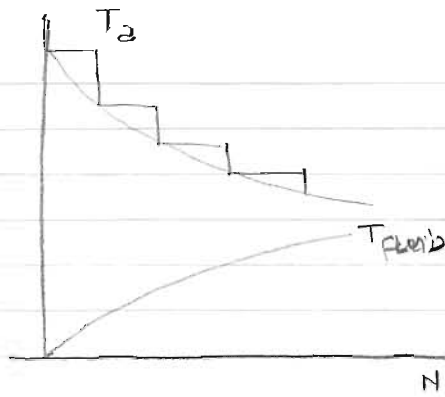
$$T_{si+1} = T_{fi} = T_{ai} - (T_{ai} - T_{si}) \exp \left(- \frac{UP}{\dot{m}c} L \right)$$

$$(T_{ai+1} - T_{si+1}) = (T_{ai} - T_{si}) \left[\exp \left(- \frac{UP}{\dot{m}c} L \right) - \frac{\dot{m} c}{\dot{m}_2 c_2} \epsilon \right]$$

$$= (T_{ai} - T_{si}) (1 - \epsilon - C \epsilon)$$

$$\boxed{\frac{Q_{i+1}}{Q_i} = 1 - (1 + C) \epsilon}$$

(d)



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$$(e) \quad Q = \sum Q_i = \dot{m} c (T_{o1} - T_{s1}) \epsilon \sum_{i=1}^{\infty} (1 - (1+c)\epsilon)^{i-1}$$

$$r = 1 - (1+c)\epsilon$$

$$\sum_{i=1}^{\infty} r^{i-1} = \frac{1}{1-r} = \frac{1}{1 - [1 - (1+c)\epsilon]} = \frac{1}{(1+c)\epsilon}$$

$$Q = \dot{m} c \frac{(T_{o1} - T_{s1})}{1+c}$$

i.e. FOR A LONG ENOUGH NUMBER OF PASSES, THE LENGTH OF THE HEAT EXCHANGER IS EQUIVALENT TO AN INFINITELY LONG EXCHANGER, WITH $\epsilon = 1$.

A COUNTER-FLOW SYSTEM HAS A HIGHER ΔT_{LMD} AND THEREFORE WOULD YIELD HIGHER HEAT TRANSFER RATES PER UNIT LENGTH.

(a) INSIDE THE ROD, WE HAVE CONDUCTION AND GENERATION:

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(i) $\nabla^2 T = \dot{q}$

$$\frac{1}{r} \frac{\partial}{\partial r} \left(\lambda r \frac{\partial T}{\partial r} \right) = \dot{q}$$

$$\lambda r \frac{\partial T}{\partial r} = \dot{q} \frac{r^2}{2} + C_1$$

$$\frac{\partial T}{\partial r} = \frac{\dot{q}}{\lambda} \frac{r}{2} + \frac{C_1}{\lambda r}$$

$$T = \frac{\dot{q}}{\lambda} \frac{r^2}{4} + \frac{C_1}{\lambda} \ln r$$

$C_1 = 0$ for finite sol.

$$T(r) = T(a) + \frac{\dot{q}}{4} \left(1 - \left(\frac{r}{a} \right)^2 \right) \quad T_c = T(0).$$

(ii) ENCLAVING TUBE: conduction only

$$\frac{1}{r} \frac{\partial}{\partial r} \left(\lambda_2 r \frac{\partial T}{\partial r} \right) = 0$$

$$\lambda_2 r \frac{\partial T}{\partial r} = C_2$$

$$\frac{\partial T}{\partial r} = \frac{C_2}{\lambda_2 r}$$

$$T = \frac{C_2}{\lambda_2} \ln r + C_3$$

$$T(a) = \frac{C_2}{\lambda_2} \ln a + C_3$$

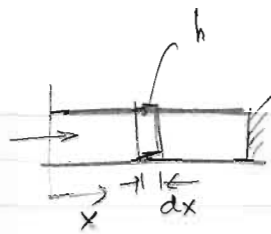
$$T(b) = \frac{C_2}{\lambda_2} \ln b + C_3$$

$$C_2 = \frac{T(b) - T(a)}{\ln(b/a)} \cdot \lambda_2$$

$$C_3 = T(a) - \frac{(T(b) - T(a)) \ln a}{\ln(b/a)}$$

$$\boxed{T(r) = T(a) + \frac{T(b) - T(a)}{\ln(b/a)} \ln \left(\frac{r}{a} \right)}$$

(b)



(i)

$$t \left\{ -\lambda \frac{\partial T}{\partial x} - \left[-\lambda \frac{\partial T}{\partial x} \right]_{x+dx} \right\} = 2h (T - T_{\infty}) dx$$

$$t \lambda \frac{\partial^2 T}{\partial x^2} = h (T - T_{\infty})$$

$$\frac{d^2 \theta}{dx^2} - \underbrace{\frac{2h}{\lambda t}}_{m^2} \theta = 0$$

$$\theta = T - T_{\infty}$$

$$\theta = A e^{mx} + B e^{-mx}$$

$$\theta(0) = T(b) - T_{\infty} = A + B$$

$$\left. \frac{d\theta}{dx} \right|_L = m \left[A e^{mL} - B e^{-mL} \right] = 0$$

$$B = A e^{mL} / e^{-mL}$$

$$A + A e^{2mL} = T(b) - T_{\infty}$$

$$A = \frac{T(b) - T_{\infty}}{1 + e^{2mL}}$$

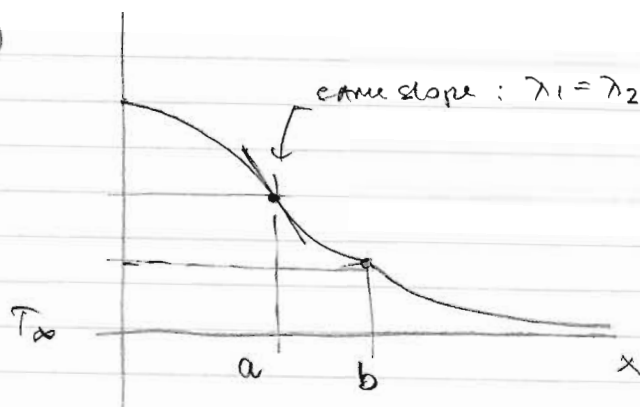
$$\therefore \theta = T - T_{\infty} = \frac{T(b) - T_{\infty}}{1 + e^{2mL}} \left(e^{mx} + e^{-mx} \frac{e^{mL}}{e^{-mL}} \right)$$

$$\left[\frac{T - T_{\infty}}{T(b) - T_{\infty}} = \frac{e^{+m(x-L)} + e^{-(x-L)}}{e^{-mL} + e^{mL}} = \frac{\cosh[m(L-x)]}{\cosh[mL]} \right]$$

(ii) h is affected by the flow pattern above the fin, whether natural or forced convection, and the properties of the fluid.

(c)

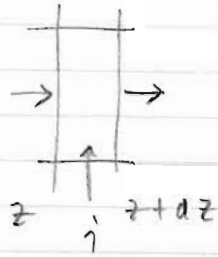
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Q2

SHB/1

(a)



$$\dot{m} \gamma_b|_z - \dot{m} \gamma_b|_{z+dz} + j_s P dz = 0$$

$$\dot{m} \gamma_{b,r} - \dot{m} \left(\gamma_{b,r} + \frac{\partial \gamma_b}{\partial z} dz \right) + j_s 2\pi R dz = 0$$

$$j_s = \frac{\dot{m}}{2\pi R} \frac{d\gamma_b}{dz} = \frac{\rho u_b}{2} \frac{d\gamma_b}{dz}$$

SINCE $\frac{\partial \gamma}{\partial z} = \text{const}$, $\frac{d\gamma_b}{dz}$ is constant and so is j_s .

$$(b) \quad u \frac{\partial \gamma}{\partial z} + \underbrace{r \frac{\partial \gamma}{\partial r}}_{\text{const.}} = D \left[\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \gamma}{\partial r} \right) \right] + \frac{\partial^2 \gamma}{\partial z^2}$$

$$u \frac{\partial \gamma}{\partial z} = \frac{D}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \gamma}{\partial r} \right)$$

$$2u_b \left(1 - \left(\frac{r}{R} \right)^2 \right) \frac{\partial \gamma}{\partial z} = \frac{D}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \gamma}{\partial r} \right)$$

SINCE $\frac{\partial \gamma}{\partial z}$ is constant, we can integrate the equation in r :

$$C_1 + \frac{2u_b}{D} \frac{\partial \gamma}{\partial z} \left(\frac{r^2}{2} - \frac{r^4}{4R^2} \right) = r \frac{\partial \gamma}{\partial r}$$

$$\frac{2u_b}{D} \frac{\partial \gamma}{\partial z} \left(\frac{r^2}{4} - \frac{r^4}{16R^2} \right) + C_1 \ln r + C_2 = \gamma$$

USING THE FACT THAT $\gamma(0) = \gamma_c$, WE HAVE:

$$\gamma - \gamma_c = \frac{u_b}{2D} \frac{\partial \gamma}{\partial z} r^2 \left(1 - \frac{1}{4} \left(\frac{r}{R} \right)^2 \right)$$

$$(c) \quad \dot{m} \gamma_b = \int_0^R \rho u \gamma \, 2\pi r \, dr$$

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$$\dot{m} (\gamma_b - \gamma_c) = \int_0^R \rho u (\gamma - \gamma_c) \, 2\pi r \, dr$$

$$\rho u_b \pi R^2 (\gamma_b - \gamma_c) = \rho 2u_b \int_0^R \left(1 - \left(\frac{r}{R}\right)^2\right) \frac{u_b}{2D} \frac{\partial \gamma}{\partial z} r^2 \left(1 - \frac{1}{4} \left(\frac{r}{R}\right)^2\right) 2\pi r \, dr$$

$$\gamma_b - \gamma_c = \frac{2}{\pi R^2} \frac{u_b}{2D} \frac{\partial \gamma}{\partial z} \int_0^R \left(r^2 - \frac{r^4}{R^2}\right) \left(1 - \frac{1}{4} \left(\frac{r^2}{R^2}\right)\right) 2\pi r \, dr$$

$$= \frac{2u_b}{DR^2} \frac{\partial \gamma}{\partial z} \int_0^R \left[r^3 - \frac{r^5}{R^2} - \frac{r^5}{4R^2} + \frac{r^7}{4R^4} \right] dr$$

$$= \frac{2u_b}{DR^2} \frac{\partial \gamma}{\partial z} \left[\frac{R^4}{4} - \frac{5}{4} \frac{R^6}{R^2} + \frac{R^8}{32 R^4} \right] =$$

$$\frac{24 - 20 + 3}{96} = \frac{7}{96}$$

$$\boxed{\gamma_b - \gamma_c = \frac{7}{48} \frac{u_b R^2}{D} \frac{\partial \gamma}{\partial z}}$$

$$(c) \quad j = \rho D \frac{\partial \gamma}{\partial r} \Big|_{r=R} = \rho D \frac{2u_b}{D} \frac{\partial \gamma}{\partial z} \left(\frac{r}{2} - \frac{r^3}{4R^2} \right) \Big|_{r=R} = \frac{\rho u_b}{2} \frac{\partial \gamma}{\partial z}$$

SINCE WE ARE IMPOSING A CONSTANT RATE OF CHANGE ON γ ,
THIS AUTOMATICALLY DIFFERS FROM THE RATE OF CHANGE

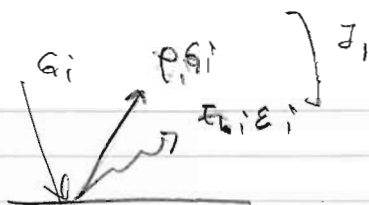
$$(d) \quad j_s = -\rho h_s (\gamma_b - \gamma_s) = -h_s [(\gamma_b - \gamma_c) - (\gamma_s - \gamma_c)] = \frac{\rho u_b}{2} \frac{\partial \gamma}{\partial z}$$

$$\gamma_s - \gamma_c = \frac{3}{8} \frac{u_b R^2}{D} \frac{\partial \gamma}{\partial z}$$

$$-h_s \left[\frac{7}{48} - \frac{3}{8} \right] \frac{u_b R^2}{D} \frac{\partial \gamma}{\partial z} = \frac{u_b}{2} \frac{\partial \gamma}{\partial z} \rightarrow h_s = \frac{48}{11} \frac{D}{2R}$$

$$\therefore Sh = \frac{h_s (2R)}{D} = \frac{48}{11}$$

Q(a)



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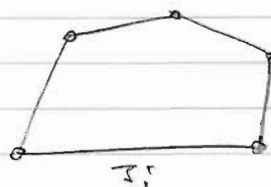
Radios. $J_i = p_i G_i + E_{b,i} \epsilon_i = (1 - \epsilon_i) G_i + E_{b,i} \epsilon_i$

$$G_i = \frac{J_i - E_{b,i} \epsilon_i}{1 - \epsilon_i}$$

Rad. Heat loss:

$$\begin{aligned} -Q_{ii} &= (J_i - G_i) A_i \\ &= A_i \left[J_i - \left(\frac{J_i - E_{b,i} \epsilon_i}{1 - \epsilon_i} \right) \right] = A_i \epsilon_i \left(\frac{E_{b,i} - J_i}{1 - \epsilon_i} \right) \quad \checkmark \end{aligned}$$

(b)



$$J_i = (1 - \epsilon_i) G_i + \epsilon_i E_{b,i}$$

$$A_i G_i = \sum_{j \neq i} F_{ji} A_j J_j + F_{ii} A_i J_i$$

$$G_i = \sum_{j \neq i} F_{ji} \frac{A_j}{A_i} J_j + F_{ii} J_i$$

$$J_i - (1 - \epsilon_i) F_{ii} J_i = (1 - \epsilon_i) \sum_{j \neq i} F_{ji} \frac{A_j}{A_i} J_j + \epsilon_i E_{b,i}$$

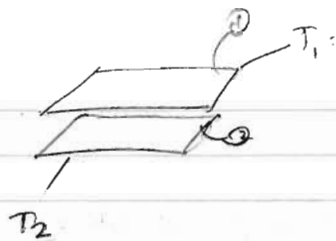
$$\text{but } F_{ji} A_j = F_{ij} A_i$$

$$\therefore J_i = \frac{1}{1 - (1 - \epsilon_i) F_{ii}} \left[(1 - \epsilon_i) \sum_{j \neq i} F_{ij} J_j + \epsilon_i E_{b,i} \right]$$

(c)

AS REQUIRED.

(c)



4/2.

③ T_3

$$F_{12} = F_{21} = 0.2$$

$$(i) \quad F_{11} + F_{12} + F_{13} = 1$$

$$F_{13} = 0.8$$

Similarly, $F_{23} = 0.8$

$$F_{13} A_1 = F_{31} A_3 \rightarrow F_{31} = 0$$

$$F_{32} = 0$$

$$F_{31} + F_{32} + F_{33} = 1 \rightarrow F_{33} = 1$$

$$\epsilon_3 = 1 \text{ (large)}$$

$$(ii) \quad J_1 = \frac{1}{1 - (1 - \epsilon_1) F_{11}} \left[(1 - \epsilon_1) F_{12} J_2 + (1 - \epsilon_1) F_{13} J_3 + \epsilon_1 E_{b1} \right]$$

$$J_2 = \frac{1}{1 - (1 - \epsilon_2) F_{22}} \left[(1 - \epsilon_2) (F_{21} J_1 + F_{23} J_3) + \epsilon_2 E_{b2} \right]$$

$$J_3 = \frac{1}{1 - (1 - \epsilon_3) F_{33}} \left[(1 - \epsilon_3) (F_{31} J_1 + F_{32} J_2) + \epsilon_3 E_{b3} \right]$$

WE HAVE :

$$J_1 - (1 - \epsilon_1) F_{12} J_2 - (1 - \epsilon_1) F_{13} J_3 = \epsilon_1 E_{b1}$$

$$-(1 - \epsilon_2) F_{21} J_1 + J_2 - (1 - \epsilon_2) F_{23} J_3 = \epsilon_2 E_{b2}$$

$$J_3 = E_{b3}$$

$$A = \begin{bmatrix} 1 & -(1 - \epsilon_1) F_{12} \\ -(1 - \epsilon_2) F_{21} & 1 \end{bmatrix} = \begin{bmatrix} 1 & -(0.2)(0.2) \\ -(0.5)(0.2) & 1 \end{bmatrix} = \begin{bmatrix} 1 & -0.04 \\ -0.1 & 1 \end{bmatrix}$$

$$C = \begin{bmatrix} \epsilon_1 E_{b1} + (1 - \epsilon_1) F_{13} E_{b3} \\ \epsilon_2 E_{b2} + (1 - \epsilon_2) F_{23} E_{b3} \end{bmatrix} = \begin{bmatrix} 0.8 E_{b1} + (0.2)(0.8) E_{b3} \\ 0.5 E_{b2} + (0.5)(0.8) E_{b3} \end{bmatrix}$$

$$\begin{bmatrix} 1 & -0.04 \\ -0.1 & 1 \end{bmatrix} \begin{bmatrix} J_1 \\ J_2 \end{bmatrix} = \begin{bmatrix} 0.8 E_{b1} + 0.16 E_{b3} \\ 0.5 E_{b2} + 0.4 E_{b3} \end{bmatrix}$$

① $\times 0.1$ + ② :

$$\begin{bmatrix} 1 & -0.04 \\ 0 & 1 - 0.004 \end{bmatrix} \begin{bmatrix} J_1 \\ J_2 \end{bmatrix} = \begin{bmatrix} 0.8 E_{b1} + 0.16 E_{b3} \\ (0.5 + 0.03) E_{b1} + (0.4 + 0.016) E_{b3} \end{bmatrix}$$

$$\underline{J_2 = 0.510 E_{b1} + 0.417 E_{b3}}$$

$$\underline{J_1 = 0.8 E_{b1} + 0.16 E_{b3} + 0.04 J_2 = 0.82 E_{b1} + 0.16 E_{b2} + 0.0167 E_{b3}}$$

$$\underline{J_3 = E_{b3}}$$

$$(iii) \quad \frac{Q}{A_1} = \frac{\epsilon_1}{1 - \epsilon_1} (E_{b1} - J_1) = \frac{0.8}{0.2} ((1 - 0.82) E_{b1} - 0.16 E_{b2} - 0.0167 E_{b3})$$

$$\underline{\frac{Q}{A_1} = 0.72 E_{b1} - 0.64 E_{b2} - 0.067 E_{b3}}$$

(d) THE HEAT LOSS BY RADIATION DOES NOT CHANGE,

BUT THE HEAT LOSS BY CONVECTION DOES.

$$\frac{Q_1}{A_1} = \frac{\epsilon_1}{1 - \epsilon_1} (E_{b1} - J_1) - h(T_1 - T_\infty)$$

THE MATHS FOR A J = C IS THE SAME, AND T_1

IS UNKNOWN. NOW WE HAVE ONE MORE EQUATION FOR

T_1 , AS ABOVE. THE NON-LINEAR SYSTEM CAN BE SOLVED

FOR T_1 AND J_1 .