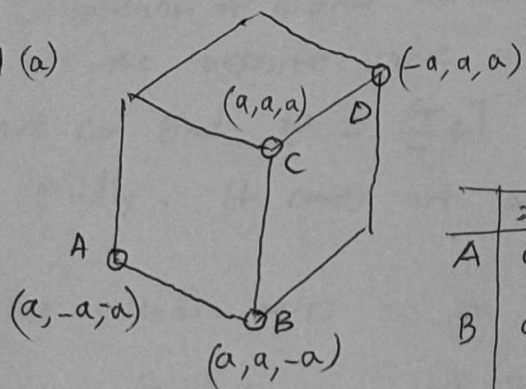


Part IIA 3C5 2026

Q1(a)



Cube $I_{xx} = I_{yy} = I_{zz} = \frac{1}{6} 2m(2a)^2$
 $= \frac{4}{3} ma^2$

$I_{xy} = I_{xz} = I_{yz} = 0$

	x_i	y_i	z_i	I_{xx}	I_{yy}	I_{zz}	I_{xy}	I_{xz}	I_{yz}
A	a	-a	-a	$2ma^2$	$2ma^2$	$2ma^2$	$-ma^2$	$-ma^2$	ma^2
B	a	a	-a	$2ma^2$	$2ma^2$	$2ma^2$	ma^2	$-ma^2$	$-ma^2$
C	a	a	a	$2ma^2$	$2ma^2$	$2ma^2$	ma^2	ma^2	ma^2
D	-a	a	a	$2ma^2$	$2ma^2$	$2ma^2$	$-ma^2$	$-ma^2$	ma^2
				$\Sigma 8ma^2$	$8ma^2$	$8ma^2$	0	$-2ma^2$	$2ma^2$

Use $I_{xx} = \Sigma m_i (y_i^2 + z_i^2)$
 $I_{yy} = \Sigma m_i (x_i^2 + z_i^2)$
 $I_{zz} = \Sigma m_i (x_i^2 + y_i^2)$

$I_{xy} = \Sigma m_i x_i y_i$, $I_{xz} = \Sigma m_i x_i z_i$, $I_{yz} = \Sigma m_i y_i z_i$

Add it all up $\therefore I_O = ma^2 \begin{bmatrix} \frac{4}{3} + 8 & 0 & 2 \\ 0 & \frac{4}{3} + 8 & -2 \\ 2 & -2 & \frac{4}{3} + 8 \end{bmatrix} = \frac{ma^2}{3} \begin{bmatrix} 28 & 0 & 6 \\ 0 & 28 & -6 \\ 6 & -6 & 28 \end{bmatrix}$

(b) $\left. \begin{aligned} \Sigma m_i x_i &= 2m \times 0 + m(a + a + a - a) = 2ma \\ \Sigma m_i y_i &= 2m \times 0 + m(-a + a + a + a) = 2ma \\ \Sigma m_i z_i &= 2m \times 0 + m(-a - a + a + a) = 0 \end{aligned} \right\} \begin{aligned} \text{total mass} &= 6m \\ \therefore x_G &= \frac{a}{3} \\ y_G &= \frac{a}{3} \\ z_G &= 0 \end{aligned}$

(c) $I_O = I_G + 6m \begin{bmatrix} y^2 + z^2 & -xy & -xz \\ -xy & x^2 + z^2 & -yz \\ -xz & -yz & x^2 + y^2 \end{bmatrix}$ with $x = -\frac{a}{3}$, $y = -\frac{a}{3}$, $z = 0$

$\therefore I_G = \frac{ma^2}{3} \begin{bmatrix} 28 & 0 & 6 \\ 0 & 28 & -6 \\ 6 & -6 & 28 \end{bmatrix} - \frac{6ma^2}{9} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 1 & 0 \\ 0 & 0 & 2 \end{bmatrix} = \frac{ma^2}{3} \begin{bmatrix} 26 & 2 & 6 \\ 2 & 26 & -6 \\ 6 & -6 & 24 \end{bmatrix}$

(note that keeping track of the signs is pretty important!)

(d) The little rotor produces $\underline{h} = \underline{I} \underline{\omega} \underline{k}$

Conservation of angular momentum means that the satellite has the opposite value of $\underline{h}$, with $\underline{h} = [\underline{I}_G] \underline{\omega}_s$

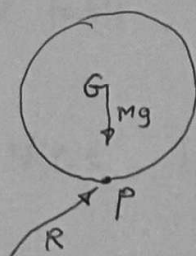
we can find $\underline{\omega}_s = [\underline{I}_G]^{-1} \begin{bmatrix} 0 \\ 0 \\ -\underline{I} \underline{\omega} \end{bmatrix}$ but it's a bit fiddly. It comes out as $\underline{\omega}_s = \frac{\underline{I} \underline{\omega}}{28 Ma^2} \begin{bmatrix} 1 \\ -1 \\ -4 \end{bmatrix}$

if the question was about verifying the result, you could have used

$$\begin{aligned} \underline{h} &= [\underline{I}_G] \underline{\omega}_s = \frac{Ma^2}{3} \begin{bmatrix} 26 & 2 & 6 \\ 2 & 26 & -6 \\ 6 & -6 & 24 \end{bmatrix} \begin{bmatrix} 1 \\ -1 \\ -4 \end{bmatrix} \times \frac{\underline{I} \underline{\omega}}{28 Ma^2} \\ &= \frac{\underline{I} \underline{\omega}}{84} \begin{bmatrix} 0 \\ 0 \\ -84 \end{bmatrix} = -\underline{I} \underline{\omega} \underline{k} \quad \checkmark \end{aligned}$$

so with this value of $\underline{\omega}_s$ the angular momentum is conserved

2 (a)



weight force passes through P

$\therefore$ net moment at P = 0

reaction at P

$$Q_P = 0$$

Use $Q_P = \underline{h}_P + \underline{r}_P \times \underline{p}$ from data sheet

with $\underline{p} = m \underline{\dot{r}}_G$

and note that P is always below G

$$\therefore \underline{r}_P = \underline{r}_G$$

$$\therefore \underline{r}_P \times \underline{p} = 0$$

$$\therefore 0 = \underline{h}_P + 0 \quad \therefore \underline{h}_P = \text{const}$$

(b) Data sheet $\underline{h}_P = \underline{h}_G + (\underline{r}_G - \underline{r}_P) \times \underline{p}$

$$\underline{p} = m(\dot{x}\underline{i} + \dot{y}\underline{j})$$

$$\underline{r}_G - \underline{r}_P = a\underline{k}$$

$$\therefore \underline{h}_P = A(\omega_1\underline{i} + \omega_2\underline{j} + \omega_3\underline{k}) + ma\underline{k} \times (\dot{x}\underline{i} + \dot{y}\underline{j})$$

with $A = \frac{2}{3}ma^2$ for a hollow sphere

$$\therefore \underline{h}_P = \left(\frac{2}{3}ma^2\omega_1 - ma\dot{y}\right)\underline{i} + \left(\frac{2}{3}ma^2\omega_2 + ma\dot{x}\right)\underline{j} + \frac{2}{3}ma^2\omega_3\underline{k}$$

= constant

$$\therefore \frac{2}{3}a\omega_1 - \dot{y} = C_1 \quad (1)$$

$$\frac{2}{3}a\omega_2 + \dot{x} = C_2 \quad (2)$$

and ω_3 is constant

(c) No slip at P means that the sphere and the table have the same velocity at P

For the table, $\underline{v}_p = \underline{\omega} \times (\underline{x} \underline{i} + \underline{y} \underline{j})$
 $= \underline{\omega} (\underline{x} \underline{j} - \underline{y} \underline{i})$ (A)

For the sphere, $\underline{v}_p = \underline{v}_G + \underline{\omega} \times \underline{r}_{p/G}$
 $= \dot{x} \underline{i} + \dot{y} \underline{j} + (\omega_1 \underline{i} + \omega_2 \underline{j} + \omega_3 \underline{k}) \times (-a \underline{k})$
 $= (\dot{x} - a\omega_2) \underline{i} + (\dot{y} + a\omega_1) \underline{j}$ (B)

No slip: (A) = (B) $\therefore \begin{cases} \dot{x} - a\omega_2 + \underline{\omega} \underline{y} = 0 & (3) \\ \dot{y} + a\omega_1 - \underline{\omega} \underline{x} = 0 & (4) \end{cases}$

from (3) $a\omega_2 = \dot{x} + \underline{\omega} \underline{y}$
 & (4) $a\omega_1 = -\dot{y} + \underline{\omega} \underline{x}$

substitute into (1) $\frac{2}{3}(-\dot{y} + \underline{\omega} \underline{x}) - \dot{y} = C_1$
 & (2) $\frac{2}{3}(\dot{x} + \underline{\omega} \underline{y}) + \dot{x} = C_2$ }

$\therefore \begin{cases} -5\dot{y} + 2\underline{\omega} \underline{x} = 3C_1 & (1a) \\ 5\dot{x} + 2\underline{\omega} \underline{y} = 3C_2 & (2a) \end{cases}$

Differentiate (2a) and substitute $\dot{x}$ from (2b)

$-5\ddot{y} + 2\underline{\omega} \underline{x} = 0$ with $\underline{x} = -\frac{2}{5}\underline{y} \underline{\omega} + \frac{3}{5}C_2$

$\therefore \ddot{y} + \left(\frac{2}{5}\underline{\omega}\right)^2 \underline{y} = \text{constant},$ (5)

Similarly differentiate (2b) and substitute $\dot{y}$ from (2a)

$\therefore \ddot{x} + \left(\frac{2}{5}\underline{\omega}\right)^2 \underline{x} = \text{constant}_2$ (6)

(5) gives (say) $\underline{y} = R \sin \frac{2}{5}\underline{\omega} \underline{t} + \underline{y}_0$
 and (6) gives $\underline{x} = R \cos \frac{2}{5}\underline{\omega} \underline{t} + \underline{x}_0$

note that from (1a) & (1b) $\underline{y}$ and $\underline{x}$ must be sin and cos

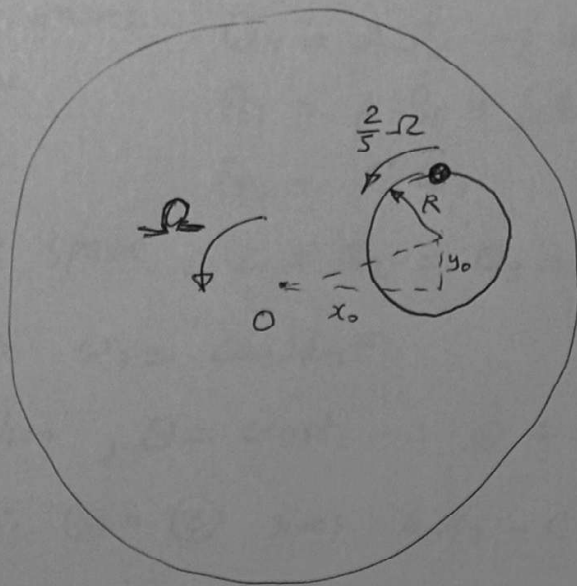
~~Now $x^2 + y^2 =$~~

$$\text{Now } (x-x_0)^2 + (y-y_0)^2 = R^2$$

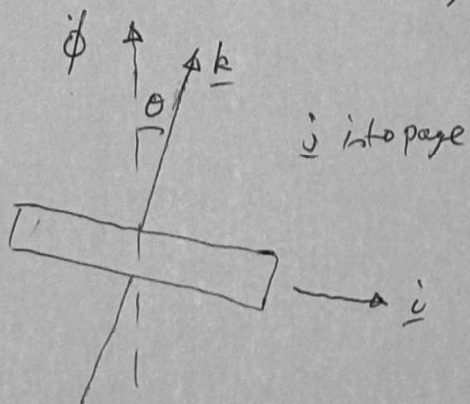
So the motion is a circle of radius R
centred on (x_0, y_0)

The frequency of the motion is $\frac{2}{5}\Omega$

So the period is $\frac{5}{2}$ times the ~~orbital~~
rotation period of the table



3(a) The triangle is an AAC body with
 $C = 2A$ by perpendicular axis theorem



Euler angles

$$\omega_1 = \Omega_1 = -\dot{\phi} \sin \theta$$

$$\omega_2 = \Omega_2 = \dot{\theta}$$

$$\Omega_3 = \dot{\phi} \cos \theta$$

$$\omega_3 = \text{spin}$$

Gyro equations

Data Sheet

$$Q_1 = A\dot{\Omega}_1 - (A\Omega_3 - C\omega_3)\Omega_2 \quad (1)$$

$$Q_2 = A\dot{\Omega}_2 + (A\Omega_3 - C\omega_3)\Omega_1 \quad (2)$$

$$Q_3 = C\dot{\omega}_3 \quad (3)$$

In free space, $Q_1 = Q_2 = Q_3 = 0$

(3) gives $\omega_3 = \text{constant}$

For nutation, $\theta = \text{const} \therefore \dot{\theta} = 0 \therefore \Omega_2 = 0$

$\therefore$ Gyro (2) gives $A\Omega_3 - C\omega_3 = 0$

$$\therefore A\dot{\phi} \cos \theta - C\omega_3 = 0$$

$$\text{For small } \theta \quad \dot{\phi} = \frac{C}{A} \omega_3 = 2\omega_3$$

This is the wobble rate and is nutation

Note that this result holds for any AAC lamina so the mass and size of the triangle don't matter

AC685, Q3b



a) $\dot{\theta}$
 $T = \frac{1}{2} m (a \dot{\theta})^2 = \frac{1}{2} J \dot{\theta}^2$
 with $J = ma^2$

$$V = m g a (1 - \cos \theta)$$

$$\approx m g a \frac{1}{2} \theta^2$$

$$L = T - V = \frac{1}{2} m a^2 \dot{\theta}^2 - m g a \frac{1}{2} \theta^2$$

b) $H = p_{\theta} \dot{\theta} - T + V$

$$p_{\theta} = \frac{\partial T}{\partial \dot{\theta}} = \frac{1}{2} m a^2 2 \dot{\theta} \Rightarrow \dot{\theta} = \frac{p_{\theta}}{m a^2}$$

$$H = \frac{p_{\theta}^2}{m a^2} - \frac{1}{2} \frac{p_{\theta}^2}{m a^2} + m g a \frac{1}{2} \theta^2$$

$$= \frac{1}{2} \frac{p_{\theta}^2}{m a^2} + m g a \frac{1}{2} \theta^2$$

c) $\frac{\partial H}{\partial p_{\theta}} = \dot{\theta} \rightarrow \frac{p_{\theta}}{m a^2} = \dot{\theta}$

$$\frac{\partial H}{\partial \theta} = -\dot{p}_{\theta} \rightarrow m g a \theta = -\dot{p}_{\theta}$$

$$\Rightarrow \frac{d}{dt} \begin{bmatrix} \theta \\ p_{\theta} \end{bmatrix} = \begin{bmatrix} \frac{p_{\theta}}{m a^2} \\ -m g a \theta \end{bmatrix}$$

d) $\theta \rightarrow x$
 $x = a \sin \theta \xrightarrow{\text{small angle}} x = a \theta$

from $(\theta, p_{\theta}) \rightarrow (x, p_x)$

$$\theta = \frac{x}{a}$$

the simplest transformation

$$G_2(q, p, t) = G_2(\theta, p_{\theta}, t) = p_x \left(\frac{x}{a} \right) = p_x \theta \Rightarrow \text{given}$$

by using the expressions in the datasheet

$$p_{\theta} = \frac{\partial G_2}{\partial \theta} = p_x$$

$$Q = \frac{\partial G_2}{\partial p_x} = x/a = \theta$$

these confirm that Hamilton's Equation with holds

Convert T in new coordinate

$$T = \frac{1}{2} m a^2 \left(\frac{\dot{x}}{a} \right)^2 = \frac{1}{2} m \dot{x}^2$$

$$p_x = \frac{\partial T}{\partial \dot{x}} = m \dot{x}$$

to find p_x

Hamiltonian in the new coordinate system

$$\Rightarrow H = K + \frac{\partial H}{\partial t} = 0 \Rightarrow K = \frac{1}{2} \frac{p_x^2}{m} + \frac{1}{2} m g a \frac{x^2}{a^2}$$

APPROACH (A)

$$H = \frac{1}{2} \frac{p_\theta^2}{ma^2} + mga \frac{1}{2} \theta^2$$

$$p_\theta = ma^2 \dot{\theta}$$

$$H = \frac{1}{2} \frac{(ma^2 \dot{\theta})^2}{ma^2} + mga \frac{1}{2} \theta^2$$

$$\theta = x/a$$

$$\Rightarrow K = \frac{1}{2} ma^2 \left(\frac{\dot{x}}{a} \right)^2 + mga \frac{1}{2} \left(\frac{x}{a} \right)^2$$

$$K = \frac{1}{2} m \dot{x}^2 + mg a \frac{1}{2} \left(\frac{x}{a} \right)^2$$

$$\Rightarrow K = \frac{1}{2} \frac{p_x^2}{m} + mga \frac{1}{2} \left(\frac{x}{a} \right)^2$$

APPROACH (B)

or find T in new coordinate

$$T = \frac{1}{2} M \dot{x}^2$$

$$V = \frac{1}{2} mag \left(\frac{x}{a} \right)^2$$

$$T = \frac{1}{2} m a^2 \dot{\theta}^2 = \frac{1}{2} m a^2 \frac{\dot{x}^2}{a^2} = \frac{1}{2} m \dot{x}^2$$

$$K = T + V$$

$$p_x = m \dot{x}$$

$$K = \frac{1}{2} \frac{p_x^2}{m} + \frac{1}{2} mag \left(\frac{x}{a} \right)^2$$

APPROACH C

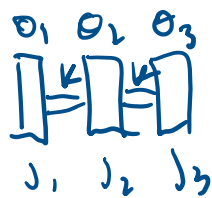
Alternatively, considering

$$p_x = p_\theta \Rightarrow ma^2 \dot{\theta} = m \dot{x} \quad \text{only true if}$$

$$\frac{1}{2} \frac{p_\theta^2}{ma^2} \Rightarrow \frac{1}{2} \frac{p_x^2}{m}$$

$$K = \frac{1}{2} \frac{p_x^2}{m} + \frac{1}{2} amg \left(\frac{x}{a} \right)^2$$

AC685, 104



a)

$$q_1 = \theta_1, \quad q_2 = \theta_2, \quad q_3 = \theta_3$$

$$T = \frac{1}{2} J_1 \dot{\theta}_1^2 + \frac{1}{2} J_2 \dot{\theta}_2^2 + \frac{1}{2} J_3 \dot{\theta}_3^2$$

$$V = \frac{1}{2} k (\theta_1 - \theta_2)^2 + \frac{1}{2} k (\theta_2 - \theta_3)^2$$

$$L = T - V$$

$$L = \frac{1}{2} J_1 \dot{\theta}_1^2 + \frac{1}{2} J_2 \dot{\theta}_2^2 + \frac{1}{2} J_3 \dot{\theta}_3^2 - \frac{1}{2} k (\theta_1 - \theta_2)^2 - \frac{1}{2} k (\theta_2 - \theta_3)^2$$

b)

$$T = \frac{1}{2} \underline{\dot{q}}^T \underline{M} \underline{\dot{q}}$$

$$\underline{M} = \begin{bmatrix} J_1 & 0 & 0 \\ 0 & J_2 & 0 \\ 0 & 0 & J_3 \end{bmatrix}$$

$$V = \frac{1}{2} \underline{q}^T \underline{K} \underline{q}$$

$$\underline{K} = \begin{bmatrix} k & -k & 0 \\ -k & 2k & -k \\ 0 & -k & k \end{bmatrix}$$

$$\begin{bmatrix} J_1 & 0 & 0 \\ 0 & J_2 & 0 \\ 0 & 0 & J_3 \end{bmatrix} \begin{bmatrix} \ddot{\theta}_1 \\ \ddot{\theta}_2 \\ \ddot{\theta}_3 \end{bmatrix} + \begin{bmatrix} c_1 & 0 & 0 \\ 0 & c_2 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \\ \dot{\theta}_3 \end{bmatrix} + \\
 + \begin{bmatrix} k & -k & 0 \\ -k & 2k & -k \\ 0 & -k & k \end{bmatrix} \begin{bmatrix} \theta_1 \\ \theta_2 \\ \theta_3 \end{bmatrix} = 0$$

$$\begin{aligned}
 c) \quad H(\theta_1, \theta_2, \theta_3, p_{\theta_1}, p_{\theta_2}, p_{\theta_3}) &= \\
 &= \sum_{j=1}^3 p_{\theta_j} \dot{\theta}_j - T + V
 \end{aligned}$$

$$p_{\theta_i} = \frac{\partial T}{\partial \dot{\theta}_i} = J_i \dot{\theta}_i \rightarrow \dot{\theta}_i = \frac{p_{\theta_i}}{J_i}$$

$$H = \frac{1}{2} \frac{p_{\theta_1}^2}{J_1} + \frac{1}{2} \frac{p_{\theta_2}^2}{J_2} + \frac{1}{2} \frac{p_{\theta_3}^2}{J_3} + \frac{1}{2} k (\theta_1 - \theta_2)^2 + \frac{1}{2} k (\theta_2 - \theta_3)^2$$

d)

$$p_{\text{TOT}} = p_1 + p_2 + p_3 = J_1 \dot{\theta}_1 + J_2 \dot{\theta}_2 + J_3 \dot{\theta}_3$$

$$\left\{ p_{\text{TOT}}, H \right\} = \underbrace{\frac{\partial p_{\text{TOT}}}{\partial \theta_1}}_{=0} \underbrace{\frac{\partial H}{\partial p_{\theta_1}}}_{=1} - \underbrace{\frac{\partial p_{\text{TOT}}}{\partial p_{\theta_1}}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_1}}_{=k(\theta_1 - \theta_2)} +$$

$$+ \underbrace{\frac{\partial p_{\text{TOT}}}{\partial \theta_2}}_{=0} \frac{\partial H}{\partial p_{\theta_2}} - \underbrace{\frac{\partial p_{\text{TOT}}}{\partial p_{\theta_2}}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_2}}_{=k(\theta_2 - \theta_3)}$$

$$+ \underbrace{\frac{\partial p_{\text{TOT}}}{\partial \theta_3}}_{=0} \frac{\partial H}{\partial p_{\theta_3}} - \underbrace{\frac{\partial p_{\text{TOT}}}{\partial p_{\theta_3}}}_{=1} \underbrace{\frac{\partial H}{\partial \theta_3}}_{=k(\theta_3 - \theta_2)}$$

$$= K(\theta_1 - \theta_2) + K(2\theta_2 - \theta_1 - \theta_3) + K(\theta_3 - \theta_2) = 0$$

However, the
total momentum is NOT conserved.

This is because there are two external generalized
damping torques applied to the system

e)

$$\frac{d}{dt} f(p, q, t) = \frac{\partial f}{\partial t} + \{f, H\}$$

$$\rightarrow \frac{dH}{dt} = \frac{\partial H}{\partial t} + \{H, H\}$$

$$\{H, H\} = 0$$

$$\text{and } \frac{\partial H}{\partial t} = 0$$

since in H
there is
NOT an explicit
time-dependence.

H is conserved.

However, this CANNOT be used as a
statement of energy conservation.

This is because of the ^{generalized} damping torques applied
on the system